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# X-RAY AND MÖSSBAUER STUDIES OF TRICYCLOHEXYLTIN(IV) HALIDES. THE CRYSTAL STRUCTURES OF <br> $\left(\text { cyclo }-\mathrm{C}_{6} \mathrm{H}_{11}\right)_{3} \operatorname{SnX}(\mathrm{X}=\mathrm{F}, \mathrm{Br}$ and I) 

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## Summary

The structures of tricyclohexyltin fluoride (I), bromide (II) and iodide (III) have been determined by X-ray analysis. Compound I crystaliizes in the space group $P 2_{1} / m$ with $a=10.422(6), b=17.238(9), c=5.769(3) \AA, \beta \doteq 104.6(1)^{\circ}$ and $Z=2$. Compounds II and III crystallize in the space group Pcmn with $a=10.427(6), b=16.914(9), c=11.366(6) \AA, Z=4$; and $a=10.400(6)$, $b=16.900(10), c=11.400(4) \AA, Z=4$, respectively. All three structures consist of discrete tetrahedral (cyclo- $\mathrm{C}_{6} \mathrm{H}_{11}$ ) 3 SnX units.

The temperature dependence of the Mössbauer resonance areas has been examined in order to obtain information about the relationship between chemical structure and lattice dynamics.

## Introduction

The tricyclohexyltin(IV) halides (cyclo- $\left.\mathrm{C}_{6} \mathrm{H}_{11}\right)_{3} \mathrm{SnX}(\mathrm{X}=\mathrm{F}, \mathrm{Cl}, \mathrm{Br}$ and I), have been subject of previous studies by Mössbauer spectroscopy [1,2]. The Mössbauer

[^0]quadrupole splitting has been interpreted in terms of a coordination geometry intermediate between four and five for the bromine and iodine compounds, and in terms of a polymeric structure involving five-coordinate tin atoms for the fluoride and chloride [2]. Recently we showed by X-ray analysis that (cyclo- $\left.\mathrm{C}_{6} \mathrm{H}_{11}\right)_{3} \mathrm{SnCl}$ consists of discrete tetrahedral units, and its quadrupole splitting value ( $\Delta=3.49 \mathrm{~mm} \mathrm{~s}^{-1}$ ) is consistent with such geometry if crystallographic data for the tin-coordination sphere are taken into account [3]. In order to gain more information about the influence of the bulky $R$ groups and the halogens on the coordination of $R_{3} \operatorname{SnX}$ compounds, the crystal structures of the title compounds have been resolved. The temperature dependence of the Mössbauer resonance areas has been also examined.

## Experimental

## Tricyclohexyltin halides

Tricyclohexyltin halides $\left(\mathrm{C}_{6} \mathrm{H}_{11}\right)_{3} \mathrm{SnX}, \mathrm{X}=\mathrm{F}, \mathrm{Br}, \mathrm{I}$, were prepared as previously described [4]. Suitable crystals for X-ray analysis were obtained from glacial acetic acid, ethanol and ethanol/petroleum ether $\left(30-60^{\circ} \mathrm{C}\right)$, respectively ( $\mathrm{X}=\mathrm{F}$, dec., $260^{\circ} \mathrm{C}$; Br, m.p. $76-77^{\circ} \mathrm{C}$; I, m.p. $63-64^{\circ} \mathrm{C}$. Lit. [4] 260, 77, $65^{\circ} \mathrm{C}$, respectively).

## Crystal and intensity data

Unit cell and space group data are listed in Table 1. Intensity data were collected by the $\theta-2 \theta$ scan method with a Philips PW 1100 computer-controlled four-circle diffractometer with graphite-monochromated Mo- $K_{\alpha}$ radiation.

As with $\left(\mathrm{C}_{6} \mathrm{H}_{11}\right)_{3} \mathrm{SnCl}$ [3], these compounds were found to decompose in the X-ray beam; the intensities fell by $6-7 \%$ of the initial values during data collection. In all cases only one crystal was used, and no absorption corrections were made in view of the relatively small values of $\mu \mathrm{R}$.

Data were collected from crystals of $\left(\text { cyclo }-\mathrm{C}_{6} \mathrm{H}_{11}\right)_{3} \mathrm{SnF}(\mathrm{I}),\left(\text { cyclo }-\mathrm{C}_{6} \mathrm{H}_{11}\right)_{3} \mathrm{SnBr}$ (II) and (cyclo- $\left.\mathrm{C}_{6} \mathrm{H}_{11}\right)_{3} \mathrm{SnI}$ (III) with dimensions $0.30 \times 0.30 \times 0.20,0.10 \times$ $0.10 \times 0.07$ and $0.20 \times 0.15 \times 0.10 \mathrm{~mm}$, respectively.

Systematic absences, and subsequent refinements, indicate space group $P 2_{1} / m$ for I and Pcmn for II and III. The intensities were corrected for Lorentz and polarization effects and scaled to give 1053, 545 and 579 independent $F_{h k l}$ values with $I>2 \sigma(I)$ for $I, \Pi$ and III, respectively.

## Solution and refinement of the structures

The structure solution was attempted for I using the coordinates of the isostructural (cyclo- $\mathrm{C}_{6} \mathrm{H}_{11}$ ) ${ }_{3} \mathrm{SnCl}$ [3], and for II and III, which are isostructural with each other, by three-dimensional Patterson methods. Full matrix least-squares refinement, with anisotropic temperature factors only for the heavy atoms; produced final $R$ factors of $0.118,0.125$ and 0.149 for $I, I I$ and III, respectively. (Hydrogen atoms were not included in the calculations.) Unitary weight was used in all calculations. Atomic scattering factors were taken from Cromer and Weber [5,6]. Final atomic coordinates and thermal parameters are listed in Table 2. The most relevant bond lengths and angles are shown in Table 3. All calculations were carried out on a CDC Cyber 76 computer using the program of ref. 7.
TABLE 1
UNIT CELL DATA

| Compound | $a(\AA)$ | $b(\AA)$ | $c$ (A) | $\beta\left({ }^{\circ}\right)$ | $U\left(\AA^{3}\right)$ | $z$ | Space group | $\begin{aligned} & D_{c} \\ & \left(\mathrm{~g} \mathrm{~cm}^{-3}\right) \end{aligned}$ | $\left(\mathrm{cm}^{-1}\right)$ | m.p. $\left({ }^{\circ} \mathrm{C}\right)$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\left(\mathrm{C}_{6} \mathrm{H}_{11}\right)_{3} \mathrm{SnF}$ | 10.422(6) | 17.238(9) | 6.769(3) | 104.6(1) | 1002.9 | 2 | $\mathrm{P}_{2} / \mathrm{m}$ | 1.28 | 12.93 |  |
| $\left(\mathrm{C}_{6} \mathrm{H}_{11}\right)_{3} \mathrm{SnCl}{ }^{\text {a }}$ | 10,386(6) | 17.195(9) | 5.713(3) | 104.3(1) | 988.5 | 2 | $\mathrm{P2}_{1} / \mathrm{m}$ | 1.35 | 14.07 | 127-128 |
| $\left(\mathrm{C}_{6} \mathrm{H}_{1}\right)_{3} \mathrm{SnBr}$ | 10.427(6) | 16.914(9) | 11.366(3) |  | 2004.5 | 4 | Pcmn | 1.48 | 33.69 | 76-77 |
| $\left(\mathrm{C}_{6} \mathrm{H}_{11}\right)_{3} \mathrm{SnI}$ | 10.400(6) | 16.900(10) | 11.400(4) |  | 2003.6 | 4 | Pcmn | 1.64 | 28.37 | 63-64 |

$a^{\text {a }}$ These data are taken from ref, 3 and repoxted for comparison.


| Com. pound | Atom | X | $\boldsymbol{Y}$ | 2 | $U_{11}$ | $U_{22}$ | $v_{33}$ | $U_{12}$ | $v_{13}$ | $U_{23}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | Sn | 359(3) | 2500 | 2494(5) | $77(2)$ | 82(2) | 98(3) | 35(6) | 31(2) | 24(6) |  |
|  | F | $331(14)$ | 2500 | -1763(25) | 59(4) |  |  |  |  |  |  |
|  | C(1) | -1760(36) | 2500 | 2228(67) | 92(11) |  |  |  |  |  |  |
|  | C(2) | -2427(56) | 3144(30) | 1022(90) | 250(24) |  |  |  |  |  |  |
|  | C(3) | -404S(64) | 3180(31) | 1241(87) | 234(22) |  |  |  |  |  |  |
|  | C(4) | -4552(60) | 2500 | -388(12) | 178(23) |  |  |  |  |  |  |
|  | C(5) | 1400(31) | 3555(17) | 3795(51) | 111 (9) |  |  |  |  |  |  |
|  | c(6) | 871 (55) | 4280(38) | 1891(101) | 246(25) |  |  |  |  |  |  |
|  | c(7) | 1688(60) | 5026(34) | 3746(96) | 248(23) |  |  |  |  |  |  |
|  | C(8) | 2933(47) | 4022(27) | 3123(78) | 181(16) |  |  |  |  |  |  |
|  | C( $\boldsymbol{\theta}^{\text {) }}$ | 3508(53) | 4376 (35) | 5305(101) | 245(29) |  |  |  |  |  |  |
|  | c(10) | 2748(51) | 3540 (31) | 4290(85) | 214(20) |  |  |  |  |  |  |
| II | Sn | 168(6) | 2500 | 681(4) | 117(5) | $79(3)$221(13) | $\begin{array}{r} 100(4) \\ 94(6) \end{array}$ | $\begin{array}{r} 40(10) \\ -61(43) \end{array}$ | $\begin{array}{r} 0(4) \\ -2(9) \end{array}$ | $\begin{gathered} 37(6) \\ 0(40) \end{gathered}$ |  |
|  | Br | 640(13) | 2500 | -1494(8) | 270(16) |  |  |  |  |  |  |
|  | c(1) | -1459(63) | 2500 | 941(48) | $77(18)$ |  |  |  |  |  |  |
|  | C(2) | -2850(90) | 1720(48) | 300(76) | 262(40) |  |  |  |  |  |  |
|  | C(3) | -4012(43) | 1791 (25) | 300(36) | 91(13) |  |  |  |  |  |  |
|  | C(4) | -4134(84) | 2500 | -291(60) | 114(22) |  |  |  |  |  |  |
|  | C(5) | 1110(73) | 1401 (44) | 1141(67) | 195(29) |  |  |  |  |  |  |
|  | C(6) | 757 (69) | 1080(60) | 2267(69) | 200(30) |  |  |  |  |  |  |
|  | C(7) | $1152(76)$ | 244 (50) | 2490(71) | 198(31) |  |  |  |  |  |  |
|  | C(8) | 2198(81) | $97(44)$ | 2002(66) | 198(31) |  |  |  |  |  |  |
|  | C(0) | 2840(80) | 430(62) | 1041(77) | 229(36) |  |  |  |  |  |  |
|  | c(10) | 2305(68) | 1320(40) | 710(57) | 172(24) |  |  |  |  |  |  |
| III | Sn | 28(7) | 2500 | 750(4) | 172(7) | $\begin{gathered} 91(4) \\ 180(15) \end{gathered}$ | $\begin{gathered} 65(3) \\ 132(10) \end{gathered}$ | $\begin{array}{r} 48(12) \\ -20(16) \end{array}$ | $\begin{gathered} 11(5) \\ 0(20) . \end{gathered}$ | $\begin{array}{r} 0(21) \\ 29(30) \end{array}$ |  |
|  | 1 | -499(7) | 2500 | -1437(6) | 154(3) |  |  |  |  |  |  |
|  | $\mathrm{C}(1)$ | -2148(68) | 2500 | $883(62)$ | $94(23)$ |  |  |  |  |  |  |
|  | c(2) | -2716(57) | 1709 (90) | 381(46) | 116(17) |  |  |  |  |  |  |
|  | c(3) | -4314(00) | 1719(34) | 446(48) | 129(20) |  |  |  |  |  |  |
|  | C(4) | -4696(73) | 2500 | $\rightarrow 93(61)$ | 126(23) |  |  |  |  |  |  |
|  | C(5) | 786 (80) | 1438(44) | 1259(70) | 166(27) |  |  |  |  |  |  |
|  | c(6) | 541(68) | 1195(45) | 2368(67) | 162(27) |  |  |  |  |  |  |
|  | c(7) | 1161(81) | 400(51) | 2786(74) | 190(33) |  |  |  |  |  |  |
|  | C(8) | 2002(89) | 38(57) | 2114(93) | 223(41) |  |  |  |  |  |  |
|  | c ${ }^{(9)}$ | 2490(83) | 339(51) | 1120(76) | 189(31) |  |  |  |  |  |  |
|  | C(10) | 1945(90) | 1213(55) | B19(75) | 204(37) |  |  |  |  |  |  |

TABLE 3
BOND LENGTHS AND BOND ANGLES INVOLVING TIN ATOM ${ }^{a}$


Tables of observed and calculated structure factors are available from the authors.

## Mössbauer measurements

Mössbauer spectra were obtained by means of the apparatus and procedures previously described [8]. The measurements, performed with a prefixed number of up-down sweeps, showed statistically undetectable Mössbauer effect above 200 K . The average amount of ${ }^{119} \mathrm{Sn}$ in the polycrystalline samples was ca. 1 mg $\mathrm{cm}^{-2}$.

The approximation relation for the high-temperature limit of the Debye model for solids, $a=\mathrm{d} \ln A / \mathrm{d} T=-1 / \mathrm{d} \mathrm{d}\left\langle x_{i}^{2}\right\rangle / \mathrm{d} T$, was employed; $a$ is the temperature coefficient, $A$ represents the Mössbauer resonance area normalized to 77.4 K , $\left\langle x_{i}^{2}\right\rangle$ is the isotropic mean-square amplitude of vibration of the tin atom and $2 \pi \lambda$ is the wavelength of the $23.9 \mathrm{keV}{ }^{119} \mathrm{Sn}$ transition [9-1.1].

As a first approximation, the isotropic mean-square amplitude of vibration of the tin atom was obtained by averaging the principal anisotropic temperature factors $U_{11}, U_{22}$ and $U_{33}$ of Table 2. These values were used to normalize to 298 K the temperature dependence of $\left\langle x_{i}^{2}\right\rangle$.

Mössbauer data for various temperature, namely the isomer shift ( $\delta$ ), quadrupole splitting ( $\Delta$ ), full-width at half-height at lower ( $\Gamma_{1}$ ) and higher energy ( $\Gamma_{2}$ ), and the logarithm of the resonance area normalized to $77.4 \mathrm{~K}(\ln A)$ are given in Table 4.

TABLE 4
MÖSSBAUER DATA FOR (cyclo- $\left.\mathrm{C}_{6} \mathrm{H}_{11}\right)_{3} \mathrm{SnX}$ COMPOUNDS

| X | $T$ (K) | $\delta a, b$ | $\Delta^{a, b}$ | $\Gamma_{1} b, c$ | $\Gamma_{2}{ }^{\text {b, }} \mathbf{c}$ | $\ln A^{d}$ | $\left(x_{i}^{2}\right)^{e}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| I | 77.4 | 1.60 | 2.76(-2.40 $f)$ | 0.78 | 0.75 | 0.000 | 0.085 |
|  | 101 | 1.60 | 2.75 | 0.73 | 0.86 | -0.380 | 0.087 |
|  | 125 | 1.58 | 2.76 | 0.85 | 0.87 | -0.766 | 0.090 |
|  | 150 | 1.55 | 2.73 | 1.03 | 0.93 | -1.168 | 0.093 |
|  | 175 | 1.59 | 2.71 | 0.83 | 0.95 | -1.570 | 0.095 |
|  | 198 | - | - | - | - | - | 0.109 |
| Br | 77.4 | 1.60 | 2.90(-2.74 $f$ ) | 0.90 | 0.93 | 0.000 | 0.074 |
|  | 101 | 1.58 | 2.89 | 0.90 | 0.86 | -0.378 | 0.077 |
|  | 125 | 1.60 | 2.91 | 0.83 | 0.92 | $-0.757$ | 0.080 |
|  | 150 | 1.59 | 2.88 | 0.90 | 0.87 | -1.190 | 0.082 |
|  | 1.75 | 1.58 | 2.90 | 0.87 | 0.86 | -1.601 | 0.085 |
|  | 298 | - | - | - | - | - | 0.099 |
| Cl | 77.4 | 1.64 | 3.50(-3.64 g) | 1.12 | 1.06 | 0.000 | 0.066 |
|  | 101 | 1.65 | 3.59 | 1.01 | 0.94 | -0.377 | 0.057 |
|  | 125 | 1.62 | 3.57 | 0.98 | 0.96 | -0.765 | 0.060 |
|  | 150 | 1.61 | 3.60 | 0.97 | 1.09 | -1.165 | 0.063 |
|  | 175 | 1.66 | 3.54 | 1.02 | 1.14 | $-1.560$ | 0.066 |
|  | 298 | - | - | - | - | - | 0.079 |
| $F$ | 77.4 | 1.55 | 3.95(-3.578) | 1.00 | 0.90 | - | - |
|  | 125 | 1.54 | 3.96 | 0.90 | 0.89 | - | - |
|  | 298 | - | - | - | - | - | 0.086 |

$a_{\text {Relative to }} \mathrm{SnO}_{2}{ }^{b} \pm \mathbf{0 . 0 5} \mathrm{mm} \mathrm{s}{ }^{-1} c_{\text {Full-widths at half-height at lower }}\left(\Gamma_{1}\right)$ and higher ( $\Gamma_{2}$ ) energy.
${ }^{1}$ Natural logarithm of the resonance area normalized to 77.4 K . $e^{\text {Amplitude of vibration of the tin }}$ atoms in $A^{2}$. $f$ Assuming an ideal tetrahedral geometry. 8 Introducing the crystallographic data in the tetrahedral geometry.

## Results and discussion

## Crystal structures

The most relevant structural features involve the coordination geometry about the tin atom. As shown in Table 3 , all the corresponding angles (except for $\mathrm{C}(1)-\mathrm{Sn}-\mathrm{Br}=111^{\circ}$ ) are almost identical within the standard deviations, including those found for $\left(\mathrm{C}_{6} \mathrm{H}_{11}\right)_{3} \mathrm{SnCl}$ [3], which has a distorted tetrahedral coordination. Thus the geometry around the tin atom is also the same in all cases.

The $\mathrm{Sn}-\mathrm{Br}(2.52 \AA$ ), $\mathrm{Sn}-\mathrm{Cl}(2.41 \AA$ ) [3] and $\mathrm{Sn}-\mathrm{C}$ (except for $\mathrm{Sn}-\mathrm{C}(1)$ of III) bond lengths all fall within the range of accepted values [12]. However the $\mathrm{Sn}-\mathrm{F}(2.45 \AA)$ and $\mathrm{Sn}-\mathrm{I}(2.54 \AA)$ bond lengths are longer and shorter, respectively, than the expected values [12]. Not too much significance should be attached to these results because of the particular experimental conditions under which the crystals were examined; these are also responsible for the relatively high final $R$ factors and for the rough geometry of the cyclohexyl groups.

The structures of $\left(\mathrm{C}_{6} \mathrm{H}_{11}\right)_{3} \mathrm{SnF}$ and of $\left(\mathrm{C}_{6} \mathrm{H}_{11}\right)_{3} \mathrm{SnCl}$ [3] consist of layers of iso-oriented molecules orthogonal to $c^{*}$, with spacing of about $5.62 \AA$. The molecules appear piled along the $c$ axis with non bonded intermolecular $\mathrm{Sn} \cdots \mathrm{F}$ distances of $3.32 \AA$ (compare $3.30 \AA$ for $\mathrm{Sn} \cdots \mathrm{Cl}$ in $\left(\mathrm{C}_{6} \mathrm{H}_{11}\right)_{3} \mathrm{SnCl}$, (see Fig. 1A).

The structures of II and III consist of layers of alternately iso- and antioriented molecules (with respect to the $a$ axis); the layers are still orthogonal to $c^{*}(\equiv c$ ) with spacings of 5.68 and $5.70 \AA$ respectively. These layers are alternately shifted by $a / 2$, giving a molecular packing in which no intermolecular Sn‥X contact distances shorter than $5 \AA$ are present (see Fig. 1B). It is noteworthy that the relatively short $\mathrm{Sn} \cdots \mathrm{X}$ contact distances in the case of the chlorine and the fluorine derivatives do not affect the main molecular geometry, which remains almost identical with that in the bromine and iodine derivatives.

## Mössbauer data

There is good agreement between our data and those previously reported [1,2]. The isomer shift values do not change within the experimental error when the halogen is varied; this reflects the constancy of the total s-electron density at the tin nucleus, as previously noted and explained for many organotin halides [13].

The following points can be made about the $\Delta$-values (cf. Table 4). The component $V_{Z Z}$ of the electric field gradient (EFG) tensor is due to a predominant valence contribution, $q_{v}$, arising from an anisotropic population in the valence shell of the $\sigma$-bonding orbitals rather than $\pi$-bondings [14], and to a small lattice contribution, $q_{1}$, arising from an anisotropic distribution of the lattice charges [15]. Therefore $q_{v}$ and $q_{1}$ have opposite signs with $\left|q_{v}\right| \gg\left|q_{i}\right|$ [16]. In our compounds the valence anisotropy arises from differences in the tin $p$-electrons involved in the $\mathrm{Sn}-\mathrm{C}$ and $\mathrm{Sn}-\mathrm{X}$ bonds. On going from the iodo to the fluoroderivative the $\Delta$-values increase with the valence anisotropy. A local pseudo-symmetry $C_{3 v}$, with the ternary axis nearly coincident with the EFG axis, may be assumed. With the further assumption that the $p$-electron density of the $\mathrm{Sn}-\mathrm{C}$ bonds is larger than that of the $\mathrm{Sn}-\mathrm{X}$ bonds, $V_{\mathrm{Zz}}$ proves to be po-

(A)

(B)

Fìg. 1. Schematic representation of the structures of the tricyclohexyltin halides as viewed along the $b$ axis. A) $X, X^{\prime}=F$ or CI. B) $X, X^{\prime}, X^{\prime \prime}=$ Br or I.
sitive for $q_{v}>0$ and $q_{1}<0 . Q$ is negative for ${ }^{119} \mathrm{Sn}$, consequently the quadrupole coupling constant, $e^{2} Q V_{\mathrm{ZZ}}$, is negative, which is in agreement with the sign of the $\Delta$-values of the cyclohexyltin halides, as calculated by the point-charge model and using the data available for tetrahedral organotin(IV) compounds [13].

Calculations of the $\Delta$-values for the cyclohexyltin halides were performed assuming the simple point-charge model and the modified model including distortions [3]: the calculated values are shown in Table 4.

The temperature coefficients $a$ derived from Mössbauer resonance areas normalized to 77.4 K (see Table 4), are ca. $-1.6 \times 10^{-2} \mathrm{~K}^{-1}$. They represent a measure of the strength of the bonds between tin and its nearest neighbour atoms [17-21]. These values are similar to those for non-interacting monomeric molecules ( $-1.8 \times 10^{-2} \mathrm{~K}^{-1}$ ) [21].

The correlation between the X-ray thermal vibration coefficients $\left\langle x_{i}^{2}\right\rangle$, for tin atoms and the temperature dependence of the Mössbauer resonance areas, $\ln A$ (see Table 4), shows that the strength of the bonds between the tin and its nearest atoms are almost constant. This appears to indicate that there are no variations in the bond strength within the series of halogen derivatives.

## Conclusion

The X-ray and Mössbauer results both indicate that tricyclohexyltin halides are characterized by discrete tetrahedral units in the crystals. The "secondary bonding" which was suggested for other similar cases [22-26] would require significantly different conformational distortions in the interacting molecules. A remarkable fact is that the molar volumes do not vary on going from fluoride to iodide derivative (see Table 1); these molar volumes seem to be mainly controlled by the large bulk of the cyclohexyl groups; and so the interactions between molecules in the crystals should be similar, in agreement with the temperature variation of the Mössbauer data.

It is noteworthy that for $\mathrm{R}_{3} \mathrm{SnX}$ compounds ( $\mathrm{X}=$ halogen) the bulky groups, such as cyclohexyl and phenyl, give rise to a tetrahedral geometry about the tin atom [3,12], whereas polymeric pentacoordinated structures are generally encountered for smaller $R$ groups [22].

The interpretation of the fact that the quadrupole splittings for the fluoride and chloride derivatives are higher than those for the bromide and iodide in terms of polymeric structures or a mixed coordination number is unsatisfactory. The structural information in the present paper indicates that the quadrupole splitting values may well be associated with differences in the crystal packing of the molecules.

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